

# Hopping Maneuver between Two Satellites: Simulations, Laboratory, and Flight Experiments on the International Space Station

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## ABSTRACT

Hopping maneuvers offer a novel approach to propellantless maneuvering of small robotic spacecraft operating in the vicinity of other spacecraft. The ongoing Astrobatics project, using NASA’s Astrobees free-flyer as a test platform, seeks to develop methodology by which the Astrobees robot can maneuver on-board the International Space Station using its robotic arm and gripper. This paper elaborates on the third experimental session of the Astrobatics project, which involves a self-toss maneuver in which a primary, active Astrobees performs a hopping maneuver off of a secondary, passive Astrobees. A MATLAB simulation of the maneuver was developed to analyze the hopping maneuver, and the results are compared with experimental maneuvers conducted on the Naval Postgraduate School’s POSEIDYN testbed, as well as at the NASA Ames Intelligent Robotics Group granite-table laboratory. From experimental results, the magnitude and direction of the launch velocity of the primary Astrobees appear to be correlated to the release angle and are consistent with simulation results.

**Keywords:** Astrobees; Astrobatics; Propellantless Self-Toss Maneuvers; Lagrangian Equations of Motion; Space Robotics; TORO

## Nomenclature

$AB_1, AB_2$	=	Astrobees Vehicle 1, Astrobees Vehicle 2
$CC_i$	=	Vectors in the inertial frame between the system CoM and the CoM of the $i^{th}$ link
$\mathbf{C}(\mathbf{x}, \dot{\mathbf{x}})$	=	Centrifugal-Coriolis Matrix
$C_i\Gamma_i$	=	Vectors between the link CoM and the augmented body CoM
$C_iA_{ij}$	=	Vectors between the link CoM and the indicated joint on that link.
$F_S$	=	Force Spring

$F_D$	= Force Damping
$G$	= Motor gear ratio
$\mathbf{H}(\mathbf{x})$	= Generalized system Inertia Matrix
$i_1 \dots i_4$	= Mass moment of Inertia (kg m <sup>2</sup> )
$K_4$	= Proportional gain
$\mathbf{L}$	= Sum of translational and rotational kinetic energies
$l_1 \dots l_4$	= Link Lengths (m)
$m_1 \dots m_4$	= Link Mass (kg)
$\mathbf{I}_{\mathcal{B}_i}^{\mathcal{C}_i}$	= $i^{\text{th}}$ link masses and inertia tensors
$\psi, \theta, \phi$	= Euler angles (heading, pitch, roll)
$\dot{\psi}, \dot{\theta}, \dot{\phi}$	= Euler angular velocity
$q_1 \dots q_6$	= Generalized Coordinates (m,m,rad,rad,rad,rad)
$\dot{q}_1 \dots \dot{q}_6$	= Generalized Coordinate Rates (m/s,m/s,rad/s,rad/s,rad/s,rad/s)
$RA$	= Release Angle
$S1 \dots S5$	= Astrobatics Sessions 1, 2, 2a, 2b, 3a, 3b, 4, and 5 respectively
$SA$	= Start Angle
$\tau$	= Manipulator Joint Torque Vector
$T_{EP}$	= kinetic energy due to the velocity of the CoM of the entire system (N)
$T_C$	= kinetic energy due to the motion of each link relative to the CoM (N)
$T$	= Total kinetic energy in the system (N)
$U$	= Potential Energy
$\mathbf{v}(\mathcal{C} \mathcal{N})$	= Velocity of the system CoM in the inertial frame
$\underline{\omega}_{\mathcal{B}_i \mathcal{N}}$	= angular velocity vector of the link bodies in the inertial frame.
$\omega_{max}$	= Maximum angular rate of the servo
$\mathbf{x} = [x_1, x_2, \dots, x_n]^T$	= Generalized manipulator joint angles

## 1 Introduction: Robotic Hopping

As space exploration advances, more national and commercial entities seek to gain access to the opportunities and resources afforded by exploring and operating beyond Earth's atmosphere. Ref. [2–4]. With this growth, the call for more efficient means of locomotion to support these operations are necessary. Of particular interest are methods of propulsion capable of reducing the need for on-board fuel, Ref. [5–10]. Conservation of momentum requires that the amount of momentum within a system remains constant, whereby a reaction force is necessary to exert a change in velocity,  $\Delta V$ , on a body.

Hopping maneuvers, as illustrated in Fig. 1, are a method of locomotion that utilize other objects, robots, or spacecraft as the reaction mass for momentum changes. This method of locomotion could be used for longer duration missions, operation of robotic assistants aboard or around large bodies such as the ISS, including autonomous vehicles that would operate on commercial satellites. Ref. [7, 11–13]

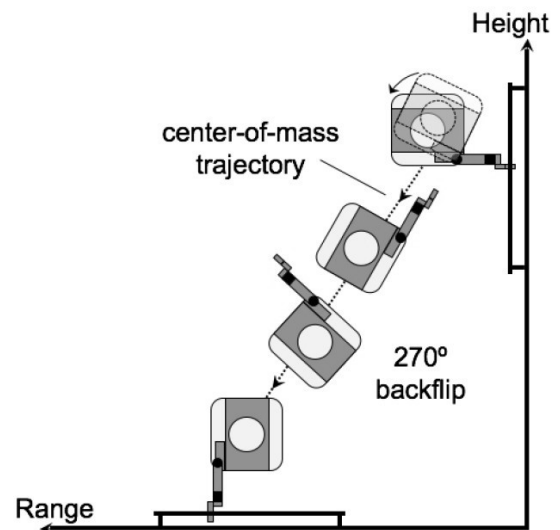
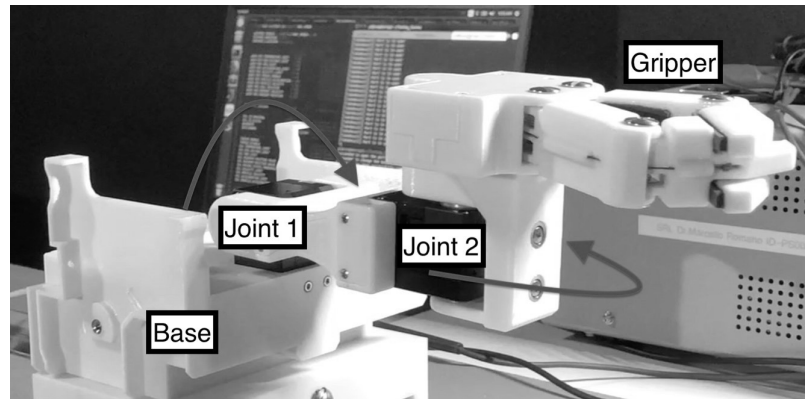


Fig. 1. Astrobatics Hop Maneuver. Ref. [1]

Astrobotics is a research collaboration between the Naval Postgraduate School (NPS) Spacecraft Robotics Laboratory (SRL) and the Intelligent Robotics Group (IRG) at NASA Ames Research Center focused on the use of the Astrobee free-flying vehicle and its robotic arm to perform self-toss maneuvers. Astrobee is equipped with a 3 degree of freedom (DOF) robotic arm composed of a two-revolute-joint, two-link manipulator and a three-finger gripper end-effector (see Fig. 2). Joint 1 and Joint 2 are respectively referred to as the proximal and distal joints of the Astrobee robotic arm. A replica of the Astrobee robotic arm was used for experimental validation tests on the NPS SRL POSEIDYN (Proximity Operation of Spacecraft: Experimental hardware-In-The-loop DYNAMIC simulator) test-bed in preparation for test sessions performed aboard the ISS. Ref. [1, 14–19].

Outlined in Tab. 1 is a summary of the five planned Astrobotics sessions. Sessions **S1** and **S2** are completed, with **S3** (the focus of this paper) having just completed. Session **S1** involved self-toss maneuvers with a single Astrobee vehicle from a handrail on the forward wall of the ISS Japanese Experimental Module (JEM), Ref. [18]. Session **S2** compared self-toss maneuvers from a deck wall handrail, a forward wall handrail, and a free-floating initial start position, followed by an impeller stabilization to bring Astrobee to rest. Session **S3** was based on the simulation and experimental validation of a type of maneuver where an active Astrobee vehicle, AB1, was commanded to perform a self-toss from a passive free-floating Astrobee vehicle, AB2. The research challenges of the **S3** session include commanding and recording data from two Astrobee vehicles simultaneously while demonstrating a robotic self-toss from a free-floating reaction mass.



**Fig. 2. NPS Replica Astrobee Arm. Joint 1 and Joint 2 are the proximal and distal joints respectively. Ref. [14]**

**Table 1 Summary of Astrobotics Sessions on-board the ISS**

Session	Date	Vehicle	Maneuver Description
<b>S1</b>	Mar. 2021	AB1	Self-toss from handrail
<b>S2a</b>	Jul. 2021	AB1	Self-toss from handrail, free-floating, with stabilization
<b>S2b</b>	Oct. 2021	AB1	
<b>S3</b>	Nov. 2021	AB1/AB2	AB1 self-toss from AB2, with stabilization
	Feb. 2022		

This paper is focused on the **S3** session of the Astrobotics project. First, an overview of Astrobee and its capabilities shall be provided with a description of the self-toss maneuvers planned to be performed during the **S3** ISS experiments. Second, a summary of the **S3** Lagrangian Equations of Motion (EOM) modeling method is included using the MATLAB TORO software developed by Professor Marcello Romano. Third, an overview of the simulation results is provided, followed by, fourth, a comparison to experiments conducted on the NPS POSEIDYN testbed and at the NASA Ames IRG granite laboratory. The testing results include descriptions of experiment performance and limitations, and data analysis of the experimental results. Finally, fifth, a summary of the first **S3** ISS session completed, with a description of conclusions and future work.

## 2 Astrobee Free-flyer Vehicle and Ground Test conditions

Astrobee first launched to the ISS in April 2019, Ref.[21], its primary function is to serve as a guest science platform for research involving free-floating robotics. Astrobee is a cube-shaped vehicle that measures 317.5 mm (12.5 in) per side and weighs 6 kg (unladen), and 10 kg (minimum-mass during operation), Ref. [20, 22]. Astrobee has a range of sensors enabling it to interact with its environment and has three payload bays that allow for guest payloads. Two impellers intake air and provide propulsion through twelve nozzles. For the Astrobotics session, Astrobee was equipped with its 3 DOF perching arm module. A diagram of the Astrobee system from the *Astrobee Guest Science Guide* is shown in Fig. 3 with further information presented in Ref. [20]. In each of the **S3** ground tests, the self-toss was initialized with two free-floating vehicles gripping the same free handrail.

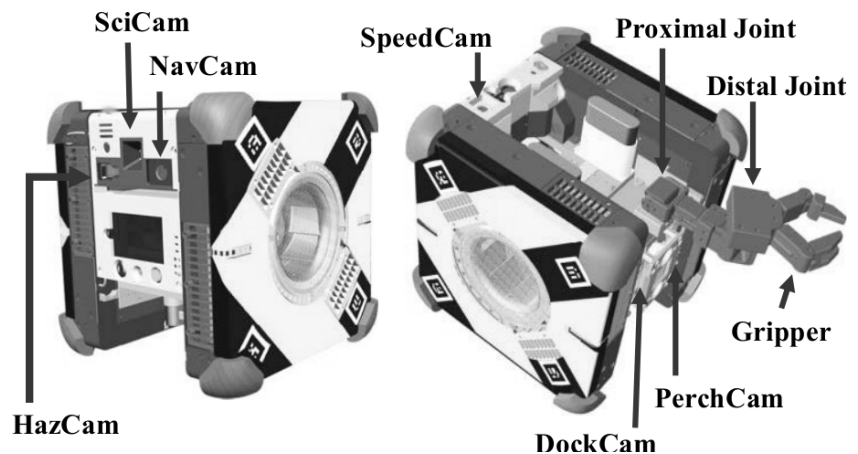


Fig. 3. NASA Astrobee Robot. Ref. [20]

The active vehicle was commanded to actuate its proximal joint from a Start Angle (SA) to a Release Angle (RA). Upon reaching the RA, the gripper end-effector was commanded to release from the handrail, causing the vehicles to move apart. The time-history of the vehicles' position and orientation after release were measured.

### 2.1 NPS POSEIDYN Test Configuration

The NPS test facility used for the Astrobotics pre-flight experimentation is the POSEIDYN testbed, which consists of a 4 m x 4 m granite monolith table and Floating Spacecraft Simulator (FSS) vehicles that operate on its surface. The POSEIDYN testbed provides the ability to conduct experiments in a simulated 3 DOF low-gravity environment Ref. [23–30].

The FSS vehicles provide mounting fixtures for a variety of experimental packages, as seen on the right side of Fig. 4. The configuration for the **S3** NPS SRL self-toss experiments is shown in Fig. 4: the active vehicle is on the left with the Astrobee arm module, and the passive vehicle is on the right.

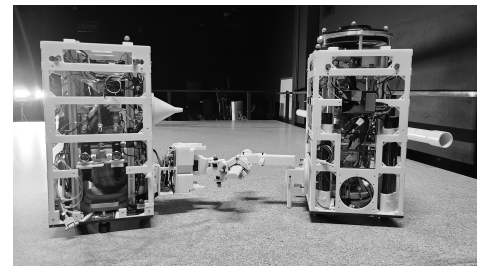


Fig. 4. Two FSS Configured for an S3 Hopping Maneuver on the POSEIDYN Testbed



Fig. 5. NASA Ames: Astrobotics S3 configuration during IRG granite lab tests

### 2.2 NASA Ames IRG Test Configuration

NASA IRG have a 2 m x 2 m granite monolith table Ref. [31]. In this facility, the Astrobee free-flyer vehicles use their native sensors to perform localization and tracking. To support testing algorithms and provide an analog of the test conditions aboard the ISS, the walls of the IRG granite lab are made to resemble the interior of the JEM Kibō



module. For the **S3** IRG test session, the Active Astrobee, AB1, was on placed on its side so that its proximal joint could move parallel to the granite surface. To make use of available mounting hardware, the Passive Astrobee, AB2, was placed in the upright orientation. A custom T-bar handrail was used as the grasping interface between both Astrobee vehicles, as shown in Fig. 5.

### 3 Lagrangian Equations of Motion

We assume the space system has zero initial linear and angular momenta; orbital effects are neglected for the short duration of the robotic maneuver. As there are no external forces acting on the system, its center of mass does not accelerate and the linear and angular momenta of the system are conserved.

The equations of motion for a free-floating space robot can be applied to the space system. Following the approach of Ref. [32], the equations of motion are

$$\mathbf{H}(\mathbf{x})\ddot{\mathbf{x}} + \mathbf{C}(\mathbf{x}, \dot{\mathbf{x}})\dot{\mathbf{x}} = \boldsymbol{\tau} \quad (1)$$

where  $\mathbf{H}(\mathbf{x})$  is the generalized system inertia matrix,  $\mathbf{C}(\mathbf{x}, \dot{\mathbf{x}})$  is the Centrifugal-Coriolis matrix, and  $\boldsymbol{\tau}$  is the manipulator joint torque vector. The manipulator joint angles form the vector  $\mathbf{x} = [x_1, x_2, \dots, x_n]^T$ .

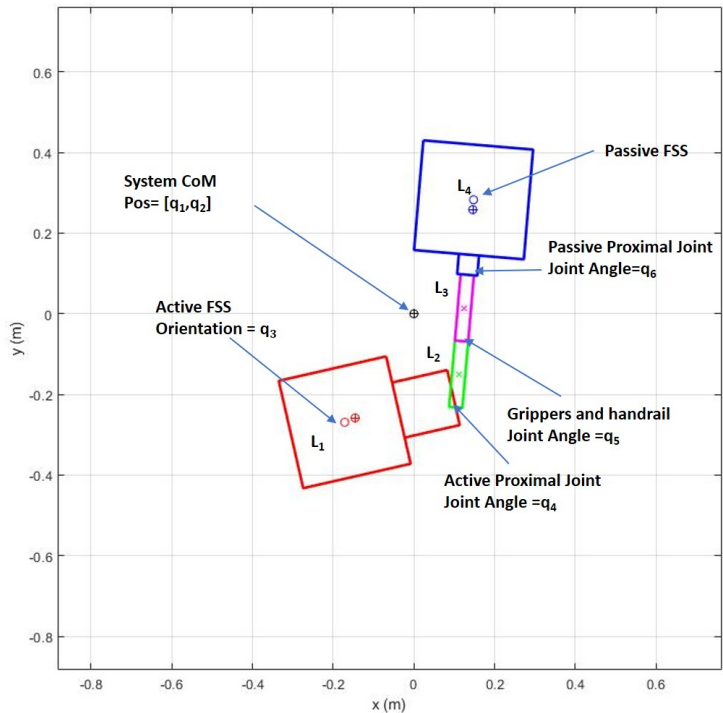
The generalized inertia matrix can be found using a Lagrangian approach,

$$\mathbf{H}(\mathbf{x}) = \frac{d}{dt} \left( \frac{d\mathbf{L}}{d\dot{\mathbf{x}}} \right) - \frac{d\mathbf{L}}{d\mathbf{x}} \quad (2)$$

Here,  $\mathbf{L}$  is the sum of the translational and rotational kinetic energies of all moving bodies in the space system; the potential energy of the system is zero.

#### 3.1 Modeling Method

The **S3** self-toss maneuver was split into a *toss* and *coast* phase. The *toss* phase began while both Astrobee vehicles gripped the handrail and was modeled as a single 4-link, 6 DOF system, whereby the dynamics of the system were propagated under the effect of AB1's joint actuation. The *toss* phase was completed when the gripper actuated and the passive Astrobee, AB2, was released from the active Astrobee, AB1. The *coast* was defined as the period after release. During the *coast*, the Astrobees were modeled as two separate, 2-link, 4 DOF systems. The two 4 DOF systems were considered free floating and maintained the rotational and linear inertia imparted during the *toss* phase with no additional internal or external forces. The gripper was assumed to instantly release and not induce any further contact forces on the system. The simulation of the *toss* and *coast* phase was implemented using the MATLAB Toolset for Orbital ROBotics (TORO) software, Ref.[33]. The virtual manipulator method, while not strictly necessary to simulate the 4-link system during the *toss*, made it possible to calculate each individual link's momentum for the *coast*.



**Fig. 6. FSS Toss Phase Diagram:** The system Center of Mass (CoM) is defined by  $q_1$  and  $q_2$ . Joint angles are given as angular displacement relative to the previous link with the exception of  $q_3$ , which is given relative to the inertial frame.

**Table 2 Symbolic Variables**

Variable Symbol	Variable Definition	Unit
$l_1 \dots l_4$	Link lengths	m
$m_1 \dots m_4$	Link Mass	kg
$ic_1 \dots ic_4$	Mass moment of inertia	$\text{kg } m^2$
$q_1 \dots q_6$	Generalized Coordinates	Various
$\dot{q}_1 \dots \dot{q}_6$	Generalized Coordinate Rates	Various

### 3.2 The Toss Phase

The kinetic and potential energy of the system was defined as a function of the generalized coordinates  $q_n$ , link dimensions ( $l_i$ ), masses ( $m_i$ ), and inertias ( $I_i$ ). Shown in Fig. 6 the red and blue squares indicate AB1 and AB2 respectively, while the green and magenta rectangles represent the Astrobee manipulators past the proximal joint. The system was oriented in the plane of motion of the Astrobee proximal joints; the distal joints were considered fixed.

The Center of Mass (CoM),  $C$ , was defined by  $q_1$  and  $q_2$  and shown as a black circle with a cross.  $q_3$  was defined as the angle of rotation of the active FSS, AB1, with respect to the inertial frame, and the coordinates  $q_4$ ,  $q_5$ , and  $q_6$  are the joint angles between the active FSS and the first arm, the first and the second arm, and the second arm and the passive FSS, respectively.

The velocity of the system CoM in the inertial frame is  $\underline{v}(C|\mathcal{N})$ , and the vectors in the inertial frame between the system CoM and the CoM of the  $i^{\text{th}}$  link were defined by  $\underline{CC}_i$ . The kinetic energy of the system in terms of the generalized coordinates and parameters is given by Eq. (3 - 5) are defined in *Orbital Space Robotics* Ref. [33]. The total kinetic energy is

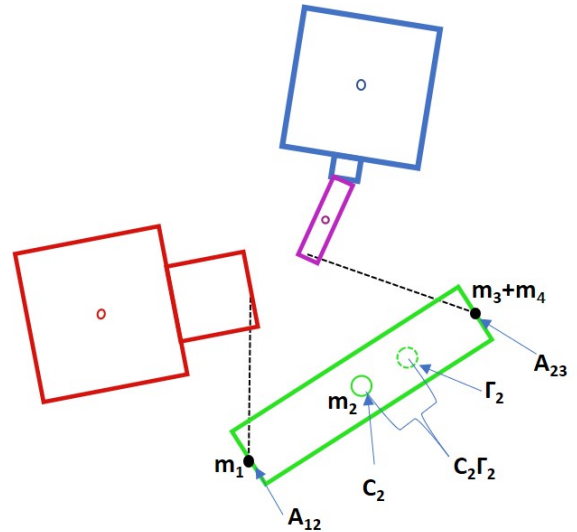
$$T = T_{EP} + T_C \quad (3)$$

where  $T_{EP}$  is defined as the kinetic energy due to the velocity of the CoM of the entire system while the term  $T_C$  is the kinetic energy due to the motion of each link relative to the CoM. Here, in addition to previously defined terms,  $m_i$  and  $\underline{\mathbf{I}}_{C_i}^{\mathcal{B}_i}$  are the  $i^{\text{th}}$  link masses and inertia tensors, and the term  $\underline{\omega}_{\mathcal{B}_i|\mathcal{N}}$  is the angular velocity vector of the link bodies in the inertial frame.

$$T_{EP} = \frac{1}{2} M \underline{v}(C|\mathcal{N}) \cdot \underline{v}(C|\mathcal{N}) \quad (4)$$

$$T_C = \frac{1}{2} \sum_{i=1}^n m_i \frac{d}{dt} \Big|_{\mathcal{N}} \underline{CC}_i \cdot \frac{d}{dt} \Big|_{\mathcal{N}} \underline{CC}_i + \frac{1}{2} \sum_{i=1}^n \underline{\omega}_{\mathcal{B}_i|\mathcal{N}} \cdot \underline{\mathbf{I}}_{C_i}^{\mathcal{B}_i} \cdot \underline{\omega}_{\mathcal{B}_i|\mathcal{N}} \quad (5)$$

To use the virtual manipulator method, the use of augmented bodies needs to be implemented. In this application, an augmented body refers to each of the links in the system with the link's CoM calculated as if the mass of the previous and proceeding attached links are point masses collocated at the link joints.



**Fig. 7. Augmented Body Illustration: Shown explicitly for link 2, the augmented body places the masses of previous and following links at the joint positions. This shifts the link's CoM to the augmented body CoM denoted by a  $\Gamma$ .**

Using Equation (6) for a four-body system [33], the vectors between the standard link CoM and the augmented body CoM were calculated for each of the bodies.  $m_i$  indicates the mass of each individual link and  $M$  indicates the total system mass. The vectors  $\underline{C}_i \underline{\Gamma}_i$  indicate the vectors between the link CoM and the augmented body CoM, while the vectors indicated by  $\underline{C}_i \underline{A}_{ij}$  are the vectors between the link CoM and the indicated joint on that link. An illustration of the second link's augmented body is shown in Fig. 7.

$$\begin{aligned} \underline{C}_1 \underline{\Gamma}_1 &= \left( \frac{m_2 + m_3 + m_4}{M} \right) \underline{C}_1 \underline{A}_{12} & \underline{C}_3 \underline{\Gamma}_3 &= \left( \frac{m_1 + m_2}{M} \right) \underline{C}_3 \underline{A}_{23} + \left( \frac{m_4}{M} \right) \underline{C}_3 \underline{A}_{34} \\ \underline{C}_2 \underline{\Gamma}_2 &= \left( \frac{m_1}{M} \right) \underline{C}_2 \underline{A}_{12} + \left( \frac{m_3 + m_4}{M} \right) \underline{C}_2 \underline{A}_{23} & \underline{C}_4 \underline{\Gamma}_4 &= \left( \frac{m_1 + m_2 + m_3}{M} \right) \underline{C}_4 \underline{A}_{34} \end{aligned} \quad (6)$$

Using the Wittenburg formula from Ref. [33] in Equation (7), the vectors indicating the displacement between the common system CoM and each individual link CoM were calculated.

$$\begin{aligned} \underline{CC}_1 &= \underline{\Gamma}_1 \underline{C}_1 + \underline{\Gamma}_2 \underline{A}_{12} + \underline{\Gamma}_3 \underline{A}_{23} + \underline{\Gamma}_4 \underline{A}_{34} & \underline{CC}_3 &= \underline{\Gamma}_1 \underline{A}_{12} + \underline{\Gamma}_2 \underline{A}_{23} + \underline{\Gamma}_3 \underline{C}_3 + \underline{\Gamma}_4 \underline{A}_{34} \\ \underline{CC}_2 &= \underline{\Gamma}_1 \underline{A}_{12} + \underline{\Gamma}_2 \underline{C}_2 + \underline{\Gamma}_3 \underline{A}_{23} + \underline{\Gamma}_4 \underline{A}_{34} & \underline{CC}_4 &= \underline{\Gamma}_1 \underline{A}_{12} + \underline{\Gamma}_2 \underline{A}_{23} + \underline{\Gamma}_3 \underline{A}_{34} + \underline{\Gamma}_4 \underline{C}_4 \end{aligned} \quad (7)$$

The time derivatives of these displacement vectors were then used to determine the kinetic energy of the system, where the angular velocities of each link relative to the inertial frame are indicated by  $\underline{\omega}_{Li\mathcal{N}}$  and iterated within TORO, Ref. [33].

$$\begin{aligned} \underline{CC}_1 &= \underline{\omega}_{L1\mathcal{N}} \times \underline{\Gamma}_1 \underline{C}_1 + \underline{\omega}_{L2\mathcal{N}} \times \underline{\Gamma}_2 \underline{A}_{12} + \underline{\omega}_{L3\mathcal{N}} \times \underline{\Gamma}_3 \underline{A}_{23} + \underline{\omega}_{L4\mathcal{N}} \times \underline{\Gamma}_4 \underline{A}_{34} \\ \underline{CC}_2 &= \underline{\omega}_{L1\mathcal{N}} \times \underline{\Gamma}_1 \underline{A}_{12} + \underline{\omega}_{L2\mathcal{N}} \times \underline{\Gamma}_2 \underline{C}_2 + \underline{\omega}_{L3\mathcal{N}} \times \underline{\Gamma}_3 \underline{A}_{23} + \underline{\omega}_{L4\mathcal{N}} \times \underline{\Gamma}_4 \underline{A}_{34} \\ \underline{CC}_3 &= \underline{\omega}_{L1\mathcal{N}} \times \underline{\Gamma}_1 \underline{A}_{12} + \underline{\omega}_{L2\mathcal{N}} \times \underline{\Gamma}_2 \underline{A}_{23} + \underline{\omega}_{L3\mathcal{N}} \times \underline{\Gamma}_3 \underline{C}_3 + \underline{\omega}_{L4\mathcal{N}} \times \underline{\Gamma}_4 \underline{A}_{34} \\ \underline{CC}_4 &= \underline{\omega}_{L1\mathcal{N}} \times \underline{\Gamma}_1 \underline{A}_{12} + \underline{\omega}_{L2\mathcal{N}} \times \underline{\Gamma}_2 \underline{A}_{23} + \underline{\omega}_{L3\mathcal{N}} \times \underline{\Gamma}_3 \underline{A}_{34} + \underline{\omega}_{L4\mathcal{N}} \times \underline{\Gamma}_4 \underline{C}_4 \end{aligned} \quad (8)$$

$$\text{where: } \underline{\omega}_{Ln\mathcal{N}} = \sum_{i=3}^{n+2} \dot{q}_i, \dot{\theta}_i = \text{norm}(\underline{\omega}_{Li\mathcal{N}}), \text{ and } : \mathbf{n}_\times = \begin{bmatrix} 0 & -1 \\ 1 & 0 \end{bmatrix} \quad (9)$$

An additional conversion was made to convert all the vectors into the same frame. The frame of the first link  $\{\hat{e}^{L1}\}$  will be used for convenience Ref. [33]. A Direction Cosine Matrix (DCM) was used to convert between the various link frames, where  $\mathbf{C}_{Li}^{Lj}$  was the DCM from the link i frame to the link j frame. For conciseness, the angular velocities were reduced to the scalar value of the z-component,  $\dot{\theta}_i$ , of the total angular velocity. Since the simulation was designed to remain in the xy-plane, the value of the z-component was determined as shown in Equation (9). The adapted equation for the time derivatives in the AB1, L1 frame is shown in Equation (10), Ref. [33].

$$\begin{aligned} [\underline{CC}_1]^{L1} &= \dot{\theta}_1 \mathbf{n}_\times [\underline{\Gamma}_1 \underline{C}_1]^{L1} + \dot{\theta}_2 \mathbf{n}_\times \mathbf{C}_{L2}^{L1} [\underline{\Gamma}_2 \underline{A}_{12}]^{L2} + \dot{\theta}_3 \mathbf{n}_\times \mathbf{C}_{L3}^{L1} [\underline{\Gamma}_3 \underline{A}_{23}]^{L3} + \dot{\theta}_4 \mathbf{n}_\times \mathbf{C}_{L4}^{L1} [\underline{\Gamma}_4 \underline{A}_{34}]^{L4} \\ [\underline{CC}_2]^{L1} &= \dot{\theta}_1 \mathbf{n}_\times [\underline{\Gamma}_1 \underline{A}_{12}]^{L1} + \dot{\theta}_2 \mathbf{n}_\times \mathbf{C}_{L2}^{L1} [\underline{\Gamma}_2 \underline{C}_2]^{L2} + \dot{\theta}_3 \mathbf{n}_\times \mathbf{C}_{L3}^{L1} [\underline{\Gamma}_3 \underline{A}_{23}]^{L3} + \dot{\theta}_4 \mathbf{n}_\times \mathbf{C}_{L4}^{L1} [\underline{\Gamma}_4 \underline{A}_{34}]^{L4} \\ [\underline{CC}_3]^{L1} &= \dot{\theta}_1 \mathbf{n}_\times [\underline{\Gamma}_1 \underline{A}_{12}]^{L1} + \dot{\theta}_2 \mathbf{n}_\times \mathbf{C}_{L2}^{L1} [\underline{\Gamma}_2 \underline{A}_{23}]^{L2} + \dot{\theta}_3 \mathbf{n}_\times \mathbf{C}_{L3}^{L1} [\underline{\Gamma}_3 \underline{C}_3]^{L3} + \dot{\theta}_4 \mathbf{n}_\times \mathbf{C}_{L4}^{L1} [\underline{\Gamma}_4 \underline{A}_{34}]^{L4} \\ [\underline{CC}_4]^{L1} &= \dot{\theta}_1 \mathbf{n}_\times [\underline{\Gamma}_1 \underline{A}_{12}]^{L1} + \dot{\theta}_2 \mathbf{n}_\times \mathbf{C}_{L2}^{L1} [\underline{\Gamma}_2 \underline{A}_{23}]^{L2} + \dot{\theta}_3 \mathbf{n}_\times \mathbf{C}_{L3}^{L1} [\underline{\Gamma}_3 \underline{A}_{34}]^{L3} + \dot{\theta}_4 \mathbf{n}_\times \mathbf{C}_{L4}^{L1} [\underline{\Gamma}_4 \underline{C}_4]^{L4} \end{aligned} \quad (10)$$

The total kinetic energy of the system was then formulated in Equation (11).

$$T_{EP} = \frac{1}{2} M (\dot{q}_1^2 + \dot{q}_2^2) \quad T_C = \frac{1}{2} \sum_{i=1}^4 m_i ([\underline{CC}_i]^{L1})' [\underline{CC}_i]^{L1} + \frac{1}{2} \sum_{i=1}^4 I_i \dot{\theta}_i^2 \quad T = T_{EP} + T_C \quad (11)$$

The system was constrained to a horizontal plane (both in the simulation and in experimentation on the granite table lab), the total potential energy,  $U$ , of the system was set as  $U = 0$ . This simulation was meant to reflect the NPS FSS, however, a small offset needs to be added to reflect the mounting hardware between the main body and the arms shown in Equation 12.

$$\begin{aligned} \underline{C_1A_{12}} &= \left[ \frac{l_1}{2} + 0.14, -0.025 \right]^\top & \underline{C_2A_{23}} &= -\underline{C_2A_{12}} & \underline{C_3A_{34}} &= -\underline{C_3A_{23}} \\ \underline{C_2A_{12}} &= \left[ -\frac{l_2}{2}, 0 \right]^\top & \underline{C_3A_{23}} &= \left[ -\frac{l_3}{2}, 0 \right]^\top & \underline{C_4A_{34}} &= \left[ -\frac{l_4}{2} - 0.05, 0 \right]^\top \end{aligned} \quad (12)$$

### 3.3 The *Coast* Phase

Developing the EOM for the *coast* phase follows a nearly identical process as the toss phase. The difference is the decrease in complexity from a 6 DOF system to two separate 4 DOF systems. Each of the two FSS two-body simulations was formulated and run separately using initial conditions provided by the toss phase simulation. Starting with formulation of the augmented body vectors in Equation 13, the formulation is immediately simpler than those in Equation 6 developed for the toss phase. This simplicity carries through to Equation 14.

$$\underline{C_1\Gamma_1} = \left( \frac{m_2}{M} \right) \underline{C_1A_{12}} \quad \underline{C_2\Gamma_2} = \left( \frac{m_1}{M} \right) \underline{C_2A_{12}} \quad (13)$$

$$\underline{C_{L2}^{L1}} = \begin{bmatrix} \cos(q_4) & -\sin(q_4) \\ \sin(q_4) & \cos(q_4) \end{bmatrix} \quad \begin{aligned} [\underline{C\dot{C}_1}]^{L1} &= \dot{\theta}_1 \mathbf{n}_\times [\underline{\Gamma_1 C_1}]^{L1} + \dot{\theta}_2 \mathbf{n}_\times \underline{C_{L2}^{L1}} [\underline{\Gamma_2 A_{12}}]^{L2} \\ [\underline{C\dot{C}_2}]^{L1} &= \dot{\theta}_1 \mathbf{n}_\times [\underline{\Gamma_1 A_{12}}]^{L1} + \dot{\theta}_2 \mathbf{n}_\times \underline{C_{L2}^{L1}} [\underline{\Gamma_2 C_2}]^{L2} \end{aligned} \quad (14)$$

The energy equations remain nearly identical to the equations used in the toss phase, but there are fewer terms in the kinetic energy equation.

$$T_{EP} = \frac{1}{2} M (\dot{q}_1^2 + \dot{q}_2^2) \quad T_C = \frac{1}{2} \sum_{i=1}^2 m_i ([\underline{C\dot{C}_i}]^{L1})' [\underline{C\dot{C}_i}]^{L1} + \frac{1}{2} \sum_{i=1}^2 I_i \dot{\theta}_i^2 \quad T = T_{EP} + T_C \quad (15)$$

### 3.4 Force Modeling

For the proximal joint, both the Astrobeer robotic arm and the NPS FSS use Dynamixel MX-106 servos. The manufacturer specifies a gear ratio of 212.6:1 [34]. Experimental data for the servo's static friction coefficient and damping ratio (0.1170 and 0.0414 respectively) are provided in Ref. [35]. The static friction and damping torque were initially set by Equations (16 - 17), where  $G$  is the gear ratio and  $\dot{q}$  is the angular rate of the joint, then further adjusted as needed.

$$F_S = 0.1179G = 25.07 \text{ Nm} \quad (16)$$

$$F_D = 0.0414\dot{q} = 8.8\dot{q} \text{ Nm} \quad (17)$$

The force applied to joint 1 (the proximal joint on the active FSS) was modelled using the advertised maximum angular rate of the servo,  $\omega_{max}$ , and a pseudo-proportional controller was developed to create a force value at each time step. The equation used for this controller is shown in Equation (18).

$$f_4 = K_4(\omega_{max} - \dot{q}_4) \quad (18)$$

The value for the proportional gain,  $K_4$ , can be set at any reasonable value that works with the simulation; however, it can be fine-tuned using experimental data to match the settling time.



## 4 Results

### 4.1 TORO S3 Simulation Results

For the NPS FSS S3 simulation, the initial parameters are listed in Table 3. The generalized and initial coordinate velocities were set to zero with the exception of  $q_4$ , which corresponds to the proximal joint position of the active FSS. The value of  $q_4$  was defined to match the starting joint angle for the desired toss maneuver. An example toss maneuver with an initial joint angle of  $90^\circ$  is shown in Tab. 3.

**Table 3 Initial Parameter Variables for Toss phase**

<b>Var.</b>	$l_1$	$l_2$	$l_3$	$l_4$	$m_1$	$m_2$	$m_3$	$m_4$	$ic_1$	$ic_2$	$ic_3$	$ic_4$
<b>Value</b>	0.273	0.165	0.165	0.273	10	1	1	10	0.124	0.002	0.002	0.124
<b>Units</b>	m	m	m	m	kg	kg	kg	kg	kg $m^2$	kg $m^2$	kg $m^2$	kg $m^2$
<b>Var.</b>	$q_1$	$q_2$	$q_3$	$q_4$	$q_5$	$q_6$	$\dot{q}_1$	$\dot{q}_2$	$\dot{q}_3$	$\dot{q}_4$	$\dot{q}_5$	$\dot{q}_6$
<b>Value</b>	0	0	0	$\frac{\pi}{2}$	0	0	0	0	0	0	0	0
<b>Units</b>	m	m	rad	rad	rad	rad	$\frac{m}{s}$	$\frac{m}{s}$	$\frac{rad}{s}$	$\frac{rad}{s}$	$\frac{rad}{s}$	$\frac{rad}{s}$

**Table 4 Initial Parameter Variables for Coast Phase from Toss Phase**

	Active Astrobee: AB1				Passive Astrobee: AB2			
<b>Var.</b>	$q_{1_{FSS1}}$	$q_{2_{FSS1}}$	$q_{3_{FSS1}}$	$q_{4_{FSS1}}$	$q_{1_{FSS2}}$	$q_{2_{FSS2}}$	$q_{3_{FSS2}}$	$q_{4_{FSS2}}$
<b>Value</b>	$\underline{R}_{FSS1}(1)$	$\underline{R}_{FSS1}(2)$	$q_3$	$q_3 + q_4$	$\underline{R}_{FSS2}(1)$	$\underline{R}_{FSS2}(2)$	$q_3 + q_4 + q_5 + q_6$	$q_6$
<b>Var.</b>	$\dot{q}_{1_{FSS1}}$	$\dot{q}_{2_{FSS1}}$	$\dot{q}_{3_{FSS1}}$	$\dot{q}_{4_{FSS1}}$	$\dot{q}_{1_{FSS2}}$	$\dot{q}_{2_{FSS2}}$	$\dot{q}_{3_{FSS2}}$	$\dot{q}_{4_{FSS2}}$
<b>Value</b>	$\dot{\underline{R}}_{FSS1}(1)$	$\dot{\underline{R}}_{FSS1}(2)$	$\dot{q}_3$	$\dot{q}_3 + \dot{q}_4$	$\dot{\underline{R}}_{FSS2}(1)$	$\dot{\underline{R}}_{FSS2}(2)$	$\dot{q}_3 + \dot{q}_4 + \dot{q}_5 + \dot{q}_6$	$\dot{q}_6$

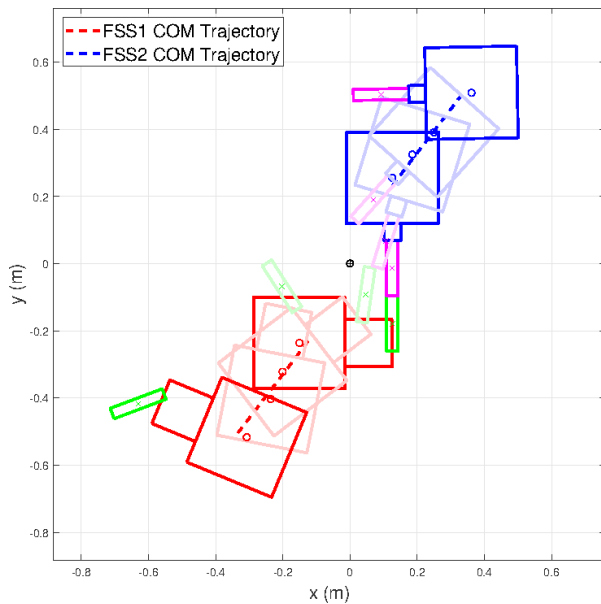
An end time,  $t_f$ , was passed to the *toss* phase simulation script, which returned a time array from  $t = 0$  to  $t = t_f$ , whose results were then used to initiate the *coast* phase. Using Equations 6 - 7 and the body vectors defined in Equation 12, the positions of each link CoM are determined relative to the system CoM.

$$\underline{C}_i = \underline{CC}_i + \begin{bmatrix} q_1 & q_2 \end{bmatrix}^\top \quad \underline{R}_{FSS1} = \frac{m_1 \underline{C}_1 + m_2 \underline{C}_2}{m_1 + m_2} \quad \underline{R}_{FSS2} = \frac{m_3 \underline{C}_3 + m_4 \underline{C}_4}{m_3 + m_4} \quad (19)$$

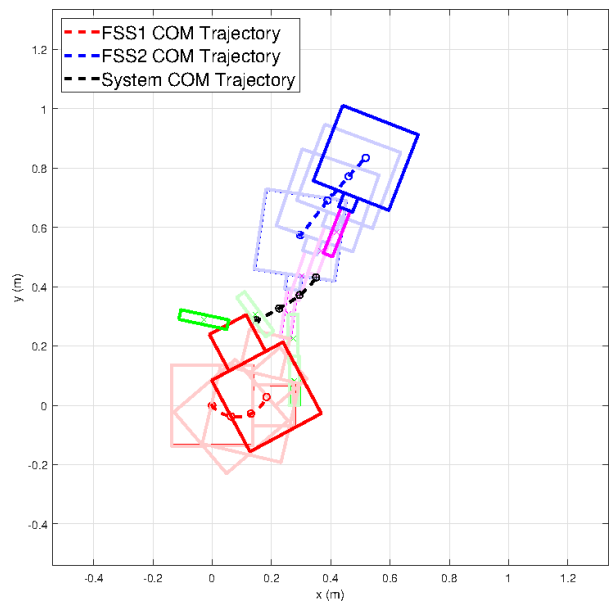
$$q_{3_{FSS1}} = q_3 \quad q_{4_{FSS1}} = q_3 + q_4 \quad q_{3_{FSS2}} = q_3 + q_4 + q_5 + q_6 \quad q_{4_{FSS2}} = q_6 \quad (20)$$

These position vectors relative to the system CoM can then be converted to the inertial frame by adding them to the position of the system CoM, which is defined by  $q_1$  and  $q_2$ . This conversion was shown in Equation 19 where  $\underline{C}_i$  is the position of link  $i$  relative to the origin.

The position of each link was converted into the positions of the CoM of each, now separate, FSS using Equation 19, where  $\underline{R}_{FSSi}$  was defined as the position vector of the  $i^{th}$  FSS CoM from the origin.  $\underline{R}_{FSSi}$  provides the first two generalized coordinates for each of the FSS. To generate the remaining generalized coordinates, the coordinates of the final state of the toss phase can be summed down the line by adding the joint angles for each link with the exception of FSS1's  $q_5$ , which is the  $q_6$  from the toss phase. Coordinate rates are generated in a similar manner by adding the joint angle rates from the end of the toss phase. The velocity of the FSS CoM was generated by taking the differences of the final two time steps and dividing them by the step size. The initial conditions provided to the two single FSS *coast* phase simulations are outlined in Tab. 4.

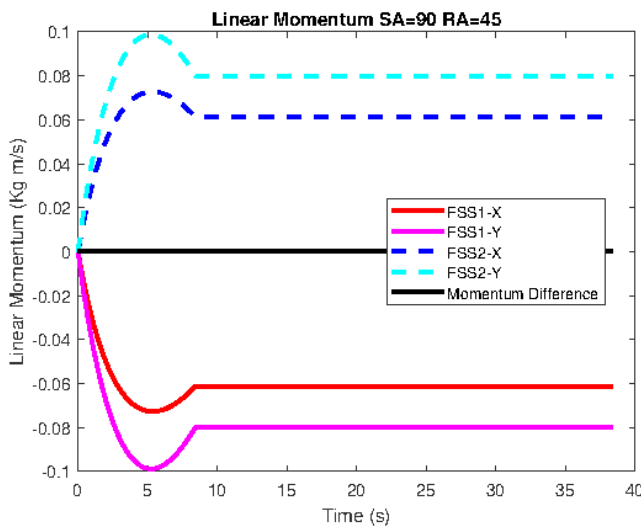


(a) SIM: Simulation Active FSS (Red) & Passive FSS (Blue). Active FSS Proximal Joint Commanded 90 to 45 Degrees

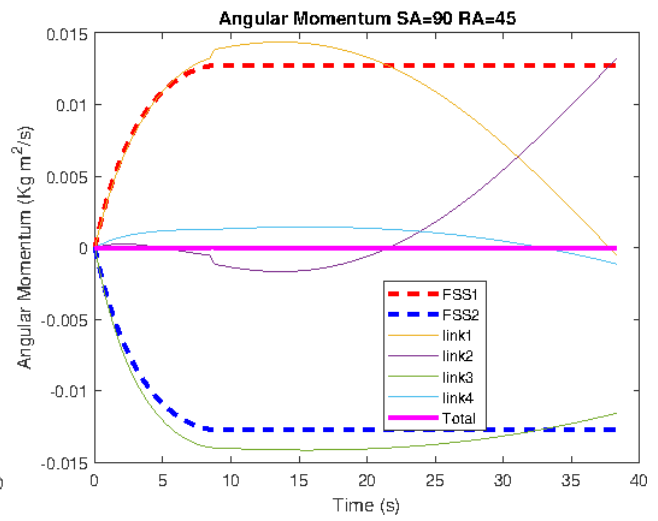


(b) EXP: NPS Experimental Self-Toss Of active FSS (red) from passive FSS (blue). Active FSS Proximal Joint commanded from SA = 90° RA = 45°

Fig. 8. Comparison Simulation to Experimental



(a) SIM: Simulation Linear Momentum



(b) SIM: Simulation Angular Momentum

Fig. 9. SIM: Linear and Angular Momentum

Using the positions and orientations over time of each link allows the entire toss maneuver to be plotted for visual reference as seen in Fig. 8a. The red squares indicate AB1 and its mounting bracket, the blue squares indicate AB2 and its mounting bracket, and the green and magenta rectangles are the corresponding arms past the proximal joint. The plot shows the initial and final positions of both FSS units as well as “ghost” positions at 10-second intervals. The dashed red and blue lines indicate the trajectories of the FSS. The stationary black circle and cross indicates that momentum was conserved during the maneuver.

This provides an immediate visual check that conservation of momentum is conserved. Fig. 9a shows a comparison of the linear momentum of each FSS in both the X and the Y directions. Here, the conservation of momentum is directly observed.

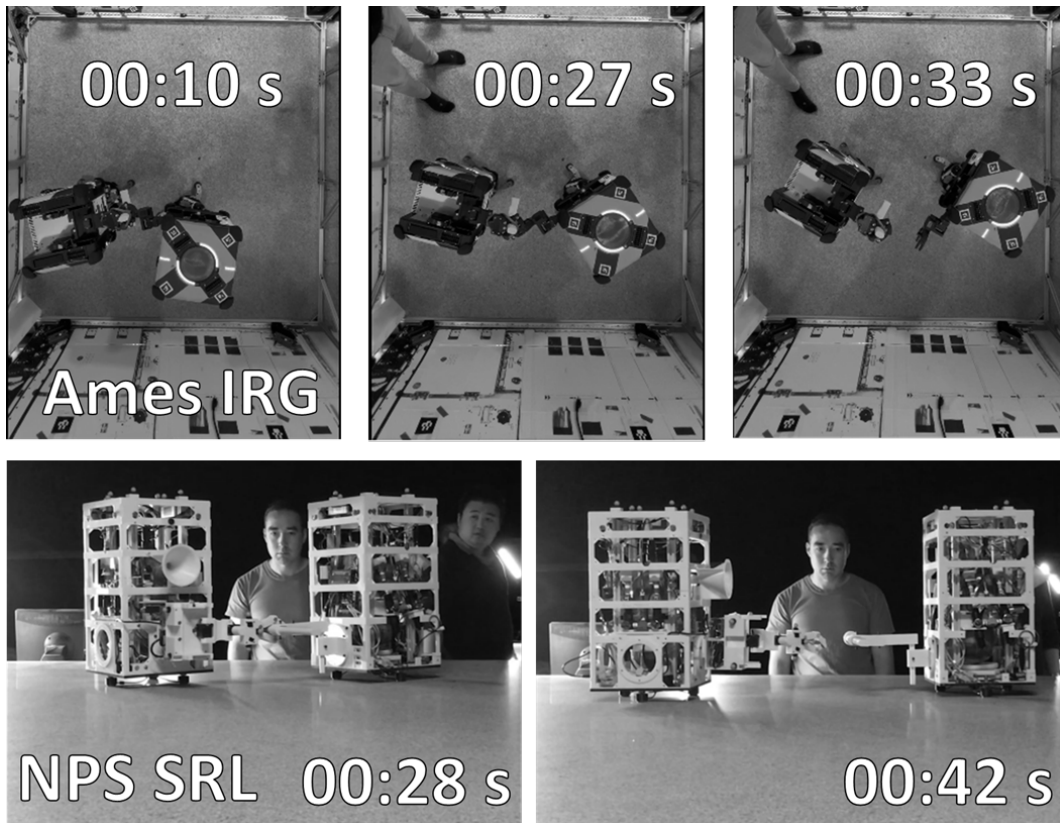


Fig. 10. SRL NPS and NASA Ames S3 Experiment

As discussed in previous section in regards to splicing the data of the *coast* and *toss* phase, the data in Fig. 9a and all future data plots include the results of both phases combined into a single continuous data array. Directly plotting the absolute value of the linear momentum and angular momenta of each individual link in Fig. 9a and 9b can further prove that momentum is conserved. The system CoM remains constant throughout the entire simulation. The CoM of each FSS move in opposite directions. Thus the data appears to show that the simulation is a physically accurate representation of a two-body self-toss maneuver.

## 4.2 NPS SRL Experimental Session

Several validation experimental runs were performed on the POSEIDYN testbed using two FSS vehicles in the configuration shown in Fig. 4, Fig. 10. Three sets of proximal toss experiments were conducted, each from a Start Angle, SA, of  $90^\circ$  to a Release Angle, RA, of  $[60, 45, 30]^\circ$  respectively. One example of the measured results is shown in Fig. 8b.

The experiments show similar self-toss trajectories as compared to the simulation; however, there appears to be a drift of the system CoM as denoted by the dashed black line in Fig. 8b. During the series of experiments, it was observed that the arm appeared to slow as it approached the RA prior to gripper release. This may have affected

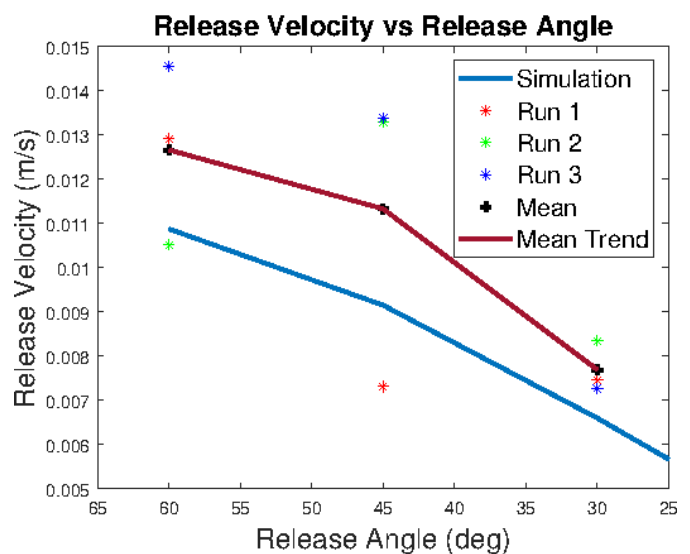


Fig. 11. EXP: Comparison of NPS Experimental Results to Simulation, where Active FSS proximal joint commanded from SA =  $90^\circ$  to RA =  $45^\circ$ : Run 1

the launch trajectory of both the active and passive FSS vehicles, causing differences between the simulations and test results.

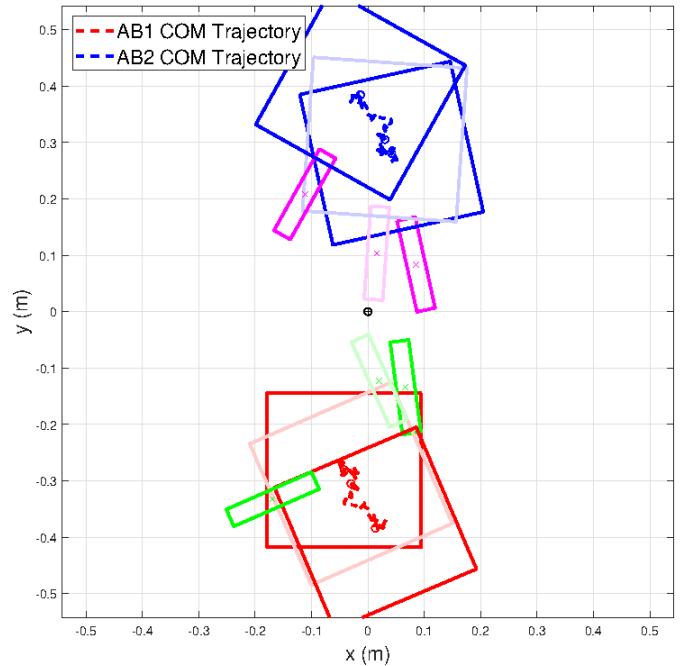
A comparison plot of the release linear velocity vs release angle for the experimental runs compared to simulation is shown in Fig. 11. The linear velocity at release for the nine experimental runs, separated by release angle, is plotted and compared to the mean simulation result found from simulation.

### 4.3 NASA Ames IRG Session

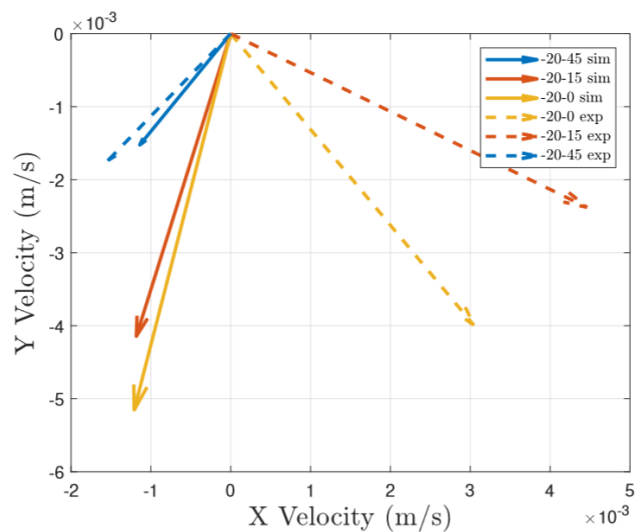
In preparation for the **S3** session aboard the ISS, a series of self-toss maneuvers were performed at the IRG granite lab. An example run where the start angle, SA, was  $-20^\circ$  to the release angle, RA, of  $0^\circ$  is shown in Fig. 12 and Fig. 10. During the test session, it was observed that there was a movement bias on the granite table, where the trajectories of both vehicles during each run moved to the top right corner of the granite table.

Possible reasons for the bias could be damaged air-bearing pads, a misalignment of the granite table, or perhaps an air pressure differential in the room. To counteract the movement bias, the relative motion of each vehicle was the plotted with respect to the system CoM as shown in Fig. 12. In so doing, the results from the POSEIDYN FSS experiments compare quite favorably to those recorded on the Ames IRG granite table. The data recorded from the Astrobees' onboard sensors is obtained by visual recognition of features in its surroundings, which differs from the external Vicon cameras used in the NPS SRL. The release velocities (both angular and linear) for each time step were calculated by dividing the displacement between the experiment start and current time elapsed. This provided an average release velocity that was then used to plot and compare the simulation vs. experimental direction vectors as shown in Fig. 13.

As shown in Fig. 13, the launch directions of the experiment results were offset from the simulation results in roughly the same direction as the offset observed during the POSEIDYN tests. The offset in the launch trajectory as compared to simulation may have been caused by similar joint actuation behaviour as observed in the NPS SRL **S3** session, here the proximal joint slowed down as it approached the desired release angle. This, coupled with possible gripper contact interaction during release, could explain the discrepancy in the simulation to experimental velocity vector for each self-toss maneuver.



**Fig. 12. EXP: Experiment at NASA Ames IRG, active (red) astrobbee proximal joint command SA =  $-20^\circ$  to RA =  $0^\circ$**



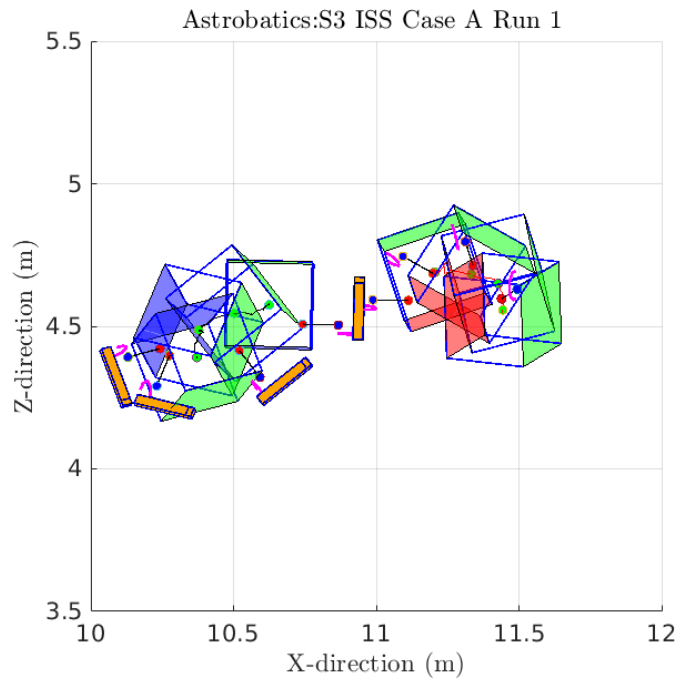
**Fig. 13. Comparison of Linear Velocity vector Directions: Simulation Vs. NASA Ames IRG Experiment**

## 4.4 ISS Test

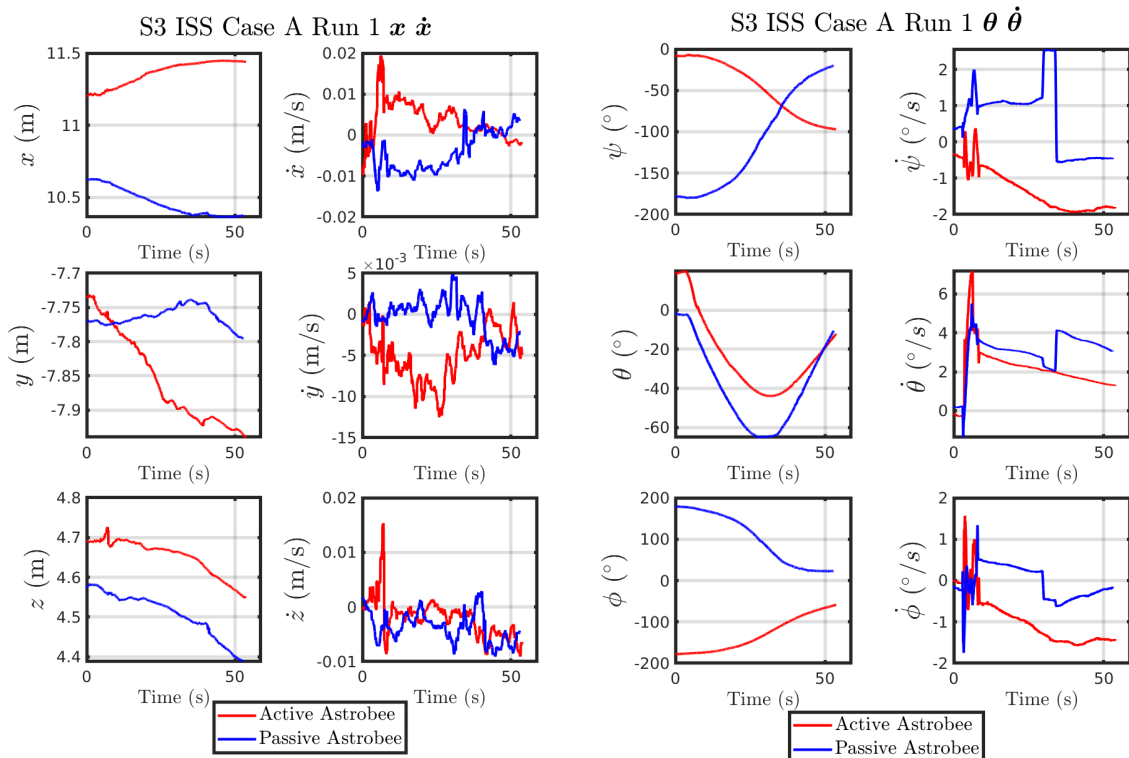
Fig. 14 is a visualization of the **S3** self-toss maneuver of the active astrobee (red) from the passive astrobee (blue). Upon self-toss, similar to the **S3** runs performed at NPS Spacecraft Robotics Laboratory and the NASA Ames Intelligent Robotics Group Astrobee testbed. The active Astrobee launches itself away from the passive Astrobee, during release both Astrobees rotate off the expected axis of actuation this could be due to the contact dynamic effects between the gripper end-effector and handrail.

Shown in Fig. 15a and Fig. 15b is the position and velocity of the active (red) and passive (blue) Astrobee vehicles during the **S3** self-toss maneuver. The self-toss and release occurs after the first 15 seconds from the start of the maneuver.

The photo in Fig. 16 was taken during the **S3** session aboard the ISS on November 23, 2021. As shown, the active Astrobee, Bumble (AB1), is on the left with the NPS skin, and the passive Astrobee, Queen (AB2), is on the right with the green skin. For each experimental run, the Astrobee vehicles were brought to rest and made to grab onto a common free-floating handrail at their initial start angles. At the start of the **S3** maneuver, the free-floating system



**Fig. 14. Time Lapse of S3 self-toss onboard the ISS**

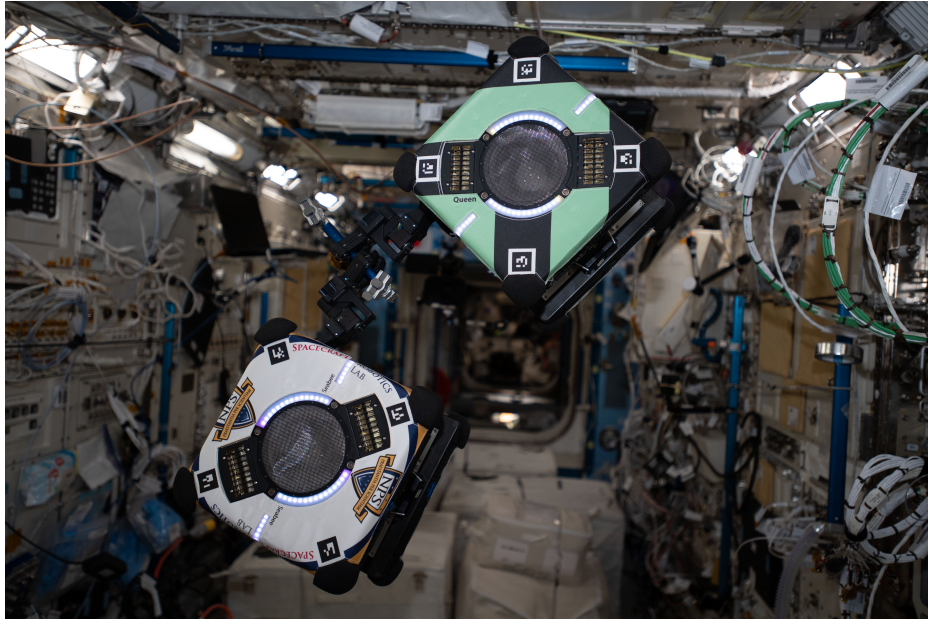


**(a) Plots of position and velocity of the active and passive astrobee vehicle during self-toss maneuver**

**(b) Plots of angular position and rate of the active and passive astrobee vehicle during self-toss maneuver**

**Fig. 15. S3 self-toss maneuver onboard the ISS**





**Fig. 16. Photo of S3 configuration during ISS session. ( Credit NASA Ref. [36])**

was released, and both AB1 and AB2 were commanded to start localization data capture, followed by AB1 performing the self-toss whilst AB2 was passively recording.

## 5 Conclusion

This paper presents a Conservation of Momentum Lagrangian method to simulate a self-toss maneuver of an active free-floating vehicle performing a self-toss maneuver from a passive free-floating object. Experimental validation was performed in the NPS SRL laboratory and at the NASA Ames IRG granite lab test facility to compare the theoretical model to the granite lab tests. This was done in preparation for an experimental session on-board the International Space Station. Overall, the simulation and granite lab tests show a correlation. However, the TORO simulation assumed perfect release of the passive object from the active vehicle with no further disturbance or contact interaction. A preliminary **S3** session of the two vehicle self-toss maneuver was conducted on-board the ISS in November 2021, as shown in Fig. 16. In addition, presented is the preliminary post-processing position and orientation results from the ISS November 2021 session.

In conclusion, a comparison of the TORO simulation to that of the experimental sessions performed NPS SRL, and Ames IRG session was able to be successfully compared with the **S3** TORO self-toss simulations, and the information gleaned can be used to predict the system response to a range of release angles.

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